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Atiyah-Singer Index Theorem Seminar.

The Heat and Wave Equations

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Heat Equation.

On IR", the heat equation is

$$\frac{\partial u_{t}}{\partial t} - \Delta u_{t}(x) = 0,$$
 (HE)

So let's think about the case on IR. Jake the Fourier transform,

$$\widehat{u}(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\infty}^{+\infty} u(x)e^{-ix\xi} dx$$

Mentle (HE) becomes

$$\frac{\partial \hat{u}}{\partial t} + \xi^2 \hat{u}_t = 0,$$

with solution $\hat{U}_{t}(\xi) = C(\xi)e^{-t\xi^{2}}$

Let V+ be such that $\frac{\partial v}{\partial v} - \Delta v_{\xi} = \bigcirc$. Then for any f, $V_{\pm} * f$ is a solution of (HE) with $V_{\circ} * f = f$. Now, $\widehat{S}_0 = \overline{\overline{\Sigma}_{000}}$, so, $\hat{V}_{\underline{t}}(\underline{\xi}) = \frac{1}{\sqrt{2\pi}} e^{-\underline{t}\underline{\xi}^2}; \quad \text{ad}$ V_E(x) = \frac{1}{\sim} \left(\frac{1}{2\pi} \left(\frac{1}{2\pi} \right) \frac{1}{2\pi} \left(\frac{1}{2\pi} \right) \frac{1}{4\pi \right} \left(\frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \left(\frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \right(\frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \frac{1}{4\pi \right} \frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} \right) \frac{1}{4\pi \right} Let's extend to Clifford Bundles ...

Let Mbe a compact Riemanian manifold, S-M a Clifford
bundle.
We have seen the motivation \mathbb{D}^2 "="- Δ .
So we define for Clifford bundles
$\frac{\partial s}{\partial t} + D^2 s = O \qquad (H.E.)$
s a section of Clifford burdle
$\frac{\partial s}{\partial t} - i \mathbb{D} s = 0$ (W.E.)
Recall that the Classical W.E. is
$\frac{\partial^2 s}{\partial t^2} - \Delta s = 0.$
OF.
Theoren:
The (WE) and (HE) have unique solutions for SoEL2(S)
L>O (HE), LER (WE)
LER (WE).
The solution satisfies: • $ s_t < s_0 $ for $t > 0$ (HE), • $ s_t = s_0 $ (WE).
• S _t = S _o (Wt)

Proof:

$$\frac{\partial}{\partial t} \|\mathbf{s}_{t}\|^{2} = \frac{\partial}{\partial t} \mathbf{s}_{t}, \mathbf{s}_{t} + \mathbf{s}_{t}, \frac{\partial}{\partial t} \mathbf{s}_{t} - \mathbf{v}^{2} \mathbf{s}_{t}, \mathbf{s}_{t} - \mathbf{s}_{t}, \mathbf{s}_{t}$$
 $= -2\|\mathbf{D}\mathbf{s}_{t}\|^{2}$
 $= -2\|\mathbf{D}\mathbf{s}_{t}\|^{2}$

heat equivariant are

Then existence follows by using the solution operator(s)

 $\mathbf{s}_{t} = e^{-tD^{2}}\mathbf{s}_{0} \quad (HE),$
 $\mathbf{s}_{t} = e^{-tD^{2}}\mathbf{s}_{0} \quad (WE)_{0}$

Since $e^{tD^{2}}$ is a smoothing operator, there exists $\mathbf{k}_{t}(\mathbf{p},\mathbf{q})$ a section of $\mathbf{S} \otimes \mathbf{S}^{*}$ such that

 $e^{-tD^{2}}\mathbf{s}_{0}(\mathbf{p}) = \int_{\mathbf{k}_{t}} \mathbf{k}_{t}(\mathbf{p},\mathbf{q}) \mathbf{s}(\mathbf{q}) \operatorname{vol}(\mathbf{q})$

and if $\mathbf{s}_{0} \in H_{t}$ then $\mathbf{s}_{t} \to \mathbf{s}_{0}$ in H_{t} .

Define \mathbf{H}_{t} and wal of \mathbf{H}_{t} via \mathbf{s}_{0} ,

 $\mathbf{q} \in H_{t}$
 $\mathbf{q} \in \mathcal{Q}^{\infty}$
 $\mathbf{q} \in \mathcal{Q}^{\infty}$

The listing enough, then the detta function the in \mathbf{H}_{t} . The solution operator them takes

 $e^{-tD^{2}}\mathbf{h}_{t} \to \mathbf{H}_{t}$, $\mathbf{H}_{t} \to \mathbf{H}_{t}$.

Theorem:

The heat hernel kt(p,q) satisfies

$$(1) \left[\frac{\partial}{\partial t} + D_p^2 \right] k_t(p,q) = 0.$$

$$\frac{\mathcal{U}}{S \otimes S_q^*} \int_{a}^{a} d D_q \text{ acts on it (is the Dince operator associated with this budge)}.$$

Moreover, k_t(p,q) is the unique such section.

Po4:

Denote
$$H = \frac{2}{3\pm} + D_p^2$$
. If we can show that

$$\frac{\text{HJ}(k_{+}(p,q) s(q) vol(q))}{=} \int_{M} H_{k_{+}(p,q)} s(q) vol(q),$$

and since the integrand in uniformly differentiable w.r.t. t, we can bring of inside the integral, and it is sufficient to show that

Locally we can write Ds= \(\sigma \); \(\nabla \); and we need only show that

$$\nabla_{i} \int_{M} k_{\ell}(p,q) \leq_{o}(q) \operatorname{vol}(q) = \int_{M} (\nabla_{i} k_{\ell}(p,q)) \leq_{o}(q) \operatorname{vol}(q).$$

Again by uniform differentiability, can bring
$$\nabla$$
; hide,

$$\nabla_i \int_{K_L} (p,q) \, S_L(q) \operatorname{vol}(q) = \int_{Y_L} (k_L(p,q) \, S_L(q)) \operatorname{vol}(q).$$

Now we have to actually do a calculation. Oy:

$$\nabla_i \left(k_L(p,q) \, S_L(p,q) \, S_L(p,q) \right) \, S_L(q) = S_L(p,q) = S_L(p,q) \, S_L(p,q) \, S_L(p,q) \, S_L(p,q) \, S_L(p,q) = S_L(p,q) \, S_$$

Finite Propogation Speed. Theorem: For any $S \in C^{\infty}(S)$, the support of $e^{i + D} s$ (ies within distance |t| of supp(s). Proof: It suffices to show that 1 | St|2 is decreasing for B(M, R) geodesic nobl of m. Why? Stare at the diagram so if the fi decreasing, \$\ls_t|^2 = 0

So let's prove that the
$$\int_{-\infty}^{\infty}$$
 is decreasing. We have that $\langle i Ds_{L}, s_{L} \rangle + \langle s_{L}, i Ds_{L} \rangle = id^{\infty}\omega$, where $\omega(x) = \langle Xs_{L}, s_{L} \rangle$.

Now,

$$\frac{\partial}{\partial L} \int_{B(h, k_{L})}^{1/2} = \int_{S(h, R_{L})}^{-\langle s_{L}, s_{L} \rangle} \cdot \langle N \cdot s_{L}, s_{L} \rangle \cdot \langle N \cdot s_{L} \rangle \cdot \langle N \cdot$$

Recall the statement from Ch5 that f -> kernel(f(D)) is continuous. Our next proposition is... Proposition: The kernel of $f(\pm D)$ artside of $\Delta \subset M \times M$ (C,S $\boxtimes S^*$) tends to zero in the C^{∞} -topology. Idea of proof: We write fu = fu + f where $\hat{f}_{u,1}$ is supported on [-8,8], and $\hat{f}_{u,2} \rightarrow 0$ in $S(\mathbb{R})$.

